

Unitarizability for the Metaplectic groups of small order

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- 2 Techniques to prove unitarizability
- 3 $\widetilde{Sp}(3)$

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Formulations

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- ❖ We denote \widetilde{P}_S as the inverse image of P_S in $\widetilde{Sp}(n)$. The Levi subgroup of \widetilde{P}_S is the quotient of $\widetilde{GL}(n_1) \times \dots \times \widetilde{GL}(n_j) \times \widetilde{Sp}(n - |S|)$.

Genuine Representation

The representation of $\widetilde{GL}(k, F)$ is obtained from tensoring a representation of $GL(k, F)$ with a genuine character of $\widetilde{GL}(k, F)$ defined by:

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Therefore for any irreducible representation π of $GL(k, F)$ the representation $\widetilde{\pi}$ of $\widetilde{GL}(k, F)$ is defined as $\widetilde{\pi}(\mathfrak{g}, \epsilon) = \pi(\mathfrak{g}) \times \chi_{\psi}(\mathfrak{g}, \epsilon)$.

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- ❖ For any ρ_i as above we can define $e(\rho_i)$ to be the unitary exponent such that $\rho_i = \nu^{e(\rho_i)} \rho^u$. Where ν denotes the determinant character and ρ^u is the unitary cuspidal genuine representations.

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Tadic structure formula

❖ Let

$$R^{\text{gen}} = \bigoplus R(\widetilde{GL(n)})_{\text{gen}},$$

where $R(\cdot)_{\text{gen}}$ denotes the Grothendieck group of finite length, smooth, genuine representation of $\widetilde{GL(n)}$.

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❖ Analogously, using the Jacquet module for the maximal parabolics of $\widetilde{GL}(n)$, we can define the following comultiplication formula:

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❖ For $\pi \in R^{\text{gen}}$ and $\sigma \in R$, we have the following formula

$$\mu^*(\pi \rtimes \alpha) = M^*(\pi) \rtimes \mu^*(\alpha).$$

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- ❖ The standard way of obtaining new unitarizable representations from old ones is by parabolic induction, which preserves unitarity (unitary parabolic induction).

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- ❖ A second way of proving unitarizability is by considering families. If a continuous family of irreducible hermitian representations of a reductive groups G contains at least one unitarizable representation, then all representations in the family are unitarizable. Furthermore, if a continuous family of irreducible hermitian representations is parameterized by unbounded set of unramified parameters, then all the representations in the family are non-unitarizable.

Irreducibility for Principal series

❖ Let $\chi_{\psi, V}$ be the character of $\widetilde{GL}(n, F)$ dependent on V and Weil index γ_{ψ} , where ψ is a non-trivial additive character of F . Let $\zeta_1, \zeta_2, \dots, \zeta_n$ be the unitary characters on F^\times . Then one can show that the following induced representation of $\widetilde{Sp}(2n, F)$:

$$\Pi = \chi_{\psi, V} \zeta_1 \times \chi_{\psi, V} \zeta_2 \times \cdots \times \chi_{\psi, V} \zeta_n \rtimes \omega_0,$$

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❖ Let $\zeta_1, \zeta_2, \dots, \zeta_n$ be the characters (not necessarily unitary) of F^\times . Suppose we have the following conditions: for any $1 \leq i \leq n$: $\zeta_i \neq \nu^{\pm 1/2} \zeta$, where ζ is a quadratic character.

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$\Pi = \chi_{\psi, \nu} \zeta_1 \times \chi_{\psi, \nu} \zeta_2 \cdots \times \chi_{\psi, \nu} \zeta_n \rtimes \omega_0$ is irreducible if and only if conditions 1 and 2 are satisfied.

Hermitian representation

Let $\langle \Delta_1, \Delta_2, \dots, \Delta_k; \sigma_{\text{neg}} \rangle$ be the unique Langlands subrepresentation of the induced representation with the segmentation. So $\langle \Delta_1, \Delta_2, \dots, \Delta_k; \sigma_{\text{neg}} \rangle$ is the hermitian if and only if:

$$\langle \Delta_1, \Delta_2, \dots, \Delta_k; \sigma_{\text{neg}} \rangle = \langle \alpha \overline{\Delta_1}, \alpha \overline{\Delta_2}, \dots, \alpha \overline{\Delta_k}; \widetilde{\overline{\sigma_{\text{neg}}}} \rangle$$

So σ_{neg} needs to be Hermitian as well. Here $\alpha = \chi_{\psi, \mathbf{V}}^2$. Therefore $\alpha \overline{\chi_{\psi, \mathbf{V}}} = \chi_{\psi, \mathbf{V}}$

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- ❖ Let $\chi_i \in \widehat{F}^\times$, then for $s_1 \geq s_2 \geq s_3 \geq 0$ $\chi_\psi \chi_1 \nu^{s_1} \times \chi_\psi \chi_2 \nu^{s_2} \times \chi_\psi \chi_3 \nu^{s_3} \rtimes \omega_0$ is irreducible except for the following cases:

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Theorem

Let $\chi_1, \chi_2 \in \widehat{F^\times}$, two nonnegative $s_1 \geq s_2 \in \mathbb{R}$, and $\zeta \in \widehat{F^\times}$ with $\zeta^2 = 1_{F^\times}$. The representations $\chi_{\psi, \nu^{s_1}} \chi_1 \times \chi_{\psi, \nu^{s_2}} \chi_2 \rtimes sp_{\zeta, 1}$ and $\chi_{\psi, \nu^{s_1}} \chi_1 \times \chi_{\psi, \nu^{s_2}} \chi_2 \rtimes \omega_{\psi_a, 1}^+$ are irreducible unless

$(s_1, \chi_1) = (3/2, \zeta)$, $(1/2, \zeta_1)$, where $\zeta_1^2 = 1_{F^\times}$, $(s_2, \chi_2) = (3/2, \zeta)$, $(1/2, \zeta_2)$, where $\zeta_2^2 = 1_{F^\times}$, or $(s_1, \chi_1) = (s_2 \pm 1, \chi_2)$.

Furthermore, if $(s_i, \chi_i) \neq (3/2, \zeta)$, $(s_i, \chi_i) \neq (1/2, \zeta_i)$ for $i = 1, 2$ and $(s_1, \chi_1) \neq (s_2 \pm 1, \chi_2)$, then $\chi_{\psi, \nu^{s_1}} \chi_1 \times \chi_{\psi, \nu^{s_2}} \chi_2 \rtimes sp_{\zeta, 1}$

$$= \begin{cases} \langle \chi_{\psi, \nu^{1/2}} \zeta; \chi_{\psi, \nu} \chi_1 \times \chi_{\psi, \nu} \chi_2 \rtimes \omega_0 \rangle & \text{if } s_1 = s_2 = 0, \\ \langle \chi_{\psi, \nu^{1/2}} \zeta, \chi_{\psi, \nu^{s_1}} \chi_1; \chi_{\psi, \nu} \chi_2 \rtimes \omega_0 \rangle & \text{if } s_2 = 0 \text{ and } s_1 \leq 1/2, \\ \langle \chi_{\psi, \nu^{1/2}} \zeta, \chi_{\psi, \nu^{s_1}} \chi_1, \chi_{\psi, \nu^{s_2}} \chi_2; \omega_0 \rangle & \text{if } 0 < s_2 \leq s_1 \leq 1/2, \\ \langle \chi_{\psi, \nu^{s_1}} \chi_1, \chi_{\psi, \nu^{1/2}} \zeta; \chi_{\psi, \nu} \chi_2 \rtimes \omega_0 \rangle & \text{if } s_2 = 0 \text{ and } 1/2 \leq s_1, \\ \langle \chi_{\psi, \nu^{s_1}} \chi_1, \chi_{\psi, \nu^{1/2}} \zeta, \chi_{\psi, \nu^{s_2}} \chi_2; \omega_0 \rangle & \text{if } 0 < s_2 \leq 1/2 \leq s_1, \\ \langle \chi_{\psi, \nu^{s_1}} \chi_1, \chi_{\psi, \nu^{s_2}} \chi_2, \chi_{\psi, \nu^{1/2}} \zeta; \omega_0 \rangle & \text{if } s_2 > 1/2, \end{cases}$$

Theorem

And $\chi_{\psi, \nu}^{s_1} \chi_1 \times \chi_{\psi, \nu}^{s_2} \chi_2 \times \omega_{\psi_a, 1}^+$

$$= \begin{cases} \chi_{\psi, \nu} \chi_1 \times \chi_{\psi, \nu} \chi_2 \times \omega_{\psi_a, 1}^+ & \text{if } s_1 = s_2 = 0, \\ \left\langle \chi_{\psi, \nu}^{\nu^{s_1}} \chi_1; \chi_{\psi, \nu} \chi_2 \times \omega_{\psi_a, 1}^+ \right\rangle & \text{if } s_2 = 0 \text{ and } s_1 > 0, \\ \left\langle \chi_{\psi, \nu}^{\nu^{s_1}} \chi_1, \chi_{\psi, \nu}^{\nu^{s_2}} \chi_2; \omega_{\psi_a, 1}^+ \right\rangle & \text{if } s_2 > 0. \end{cases}$$

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 We also have the following: Let $\chi, \chi_1 \in \widehat{F^\times}$. Suppose s is non-negative, $s_1 \geq s_2$ and $s_3 \in \mathbb{R}$. Suppose $\zeta, \zeta_1 \in \widehat{F^\times}$ such that $\zeta^2 = \zeta_1^2 = 1$. Then the representations $\chi_{\psi, \nu^s} \chi \text{St}_{GL(2)} \times \chi_{\psi, \nu^{s_3}} \chi_1 \rtimes \omega_0$ and $\chi_{\psi, \nu^s} \chi \mathbb{1}_{GL(2)} \times \chi_{\psi, \nu^{s_3}} \chi_1 \rtimes \omega_0$ are irreducible except in the case where $(s_3, \chi_1) = (1/2, \zeta_1)$, $(s, \chi) = (1/2, \zeta)$ or $(1, \zeta)$ or $(0, \zeta)$, $(s_3, \chi_1) = (s \pm 1/2, \chi)$, $(s \pm 3/2, \chi)$. Also if $\chi_{\psi, \nu^{s_1}} \chi \times \chi_{\psi, \nu^{s_2}} \chi \times \chi_{\psi, \nu^{s_3}} \chi_1 \rtimes \omega_0$ is a representation of length 2, then we have $\chi_{\psi, \nu^s} \chi \mathbb{1}_{GL(2)} \times \chi_{\psi, \nu^{s_3}} \chi_1 \rtimes \omega_0$

$$= \begin{cases} \langle \chi_{\psi, \nu^s} \chi \mathbb{1}_{GL(2)}, \chi_{\psi, \nu^{s_3}} \chi_1; \omega_0 \rangle & \text{if } s_3 > 0 \\ \langle \chi_{\psi, \nu^s} \chi \mathbb{1}_{GL(2)}; \chi_{\psi, \nu} \chi_1 \rtimes \omega_0 \rangle & \text{if } s_3 = 0 \end{cases}$$

Theorem

For the other case we have $\chi_{\psi, \nu^s} \chi_{StGL(2)} \times \chi_{\psi, \nu^{s_3}} \chi_1 \rtimes \omega_0 =$

$$\left\{ \begin{array}{ll} \langle \chi_{\psi, \nu^{s+1/2}} \chi, \chi_{\psi, \nu^{s-1/2}} \chi, \chi_{\psi, \nu^{s_3}} \chi_1; \omega_0 \rangle & s - 1/2 > s_3 \\ \langle \chi_{\psi, \nu^{s+1/2}} \chi, \chi_{\psi, \nu^{1/2-s}} \chi, \chi_{\psi, \nu^{s_3}} \chi_1; \omega_0 \rangle & 1/2 - s > s_3, s < 1/2 \\ \langle \chi_{\psi, \nu^{s+1/2}} \chi, \chi_{\psi, \nu^{1/2-s}} \chi; \chi_{\psi, \nu} \chi_1 \rtimes \omega_0 \rangle & s_3 = 0, \text{ and } s < 1/2 \\ \langle \chi_{\psi, \nu} \chi; \chi_{\psi, \nu} \chi \times \chi_{\psi, \nu} \chi_1 \rtimes \omega_0 \rangle & s_3 = 0 \text{ and } s = 1/2 \\ \langle \chi_{\psi, \nu} \chi, \chi_{\psi, \nu^{s_3}} \chi_1; \chi_{\psi, \nu} \chi \rtimes \omega_0 \rangle & 1 > s_3 > 0 \text{ and } s = 1/2 \\ \langle \chi_{\psi, \nu^{s+1/2}} \chi, \chi_{\psi, \nu^{s_3}} \chi_1, \chi_{\psi, \nu^{s-1/2}} \chi; \omega_0 \rangle & s + 1/2 > s_3 > s - 1/2 \\ \langle \chi_{\psi, \nu^{s_3}} \chi, \chi_{\psi, \nu^{s+1/2}} \chi_1, \chi_{\psi, \nu^{s-1/2}} \chi; \omega_0 \rangle & s_3 > s + 1/2, s > 1/2 \\ \langle \chi_{\psi, \nu^{s_3}} \chi_1, \chi_{\psi, \nu^{s+1/2}} \chi; \chi_{\psi, \nu} \chi \rtimes \omega_0 \rangle & s_3 > s + 1/2, s = 1/2 \\ \langle \chi_{\psi, \nu^{s_3}} \chi, \chi_{\psi, \nu^{s+1/2}} \chi_1, \chi_{\psi, \nu^{1/2-s}} \chi; \omega_0 \rangle & s_3 > s + 1/2, s < 1/2 \end{array} \right.$$

Exceptional cases

They are listed as below:

$$\chi_{\psi, \nu} \zeta^{\nu^{3/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \omega_0.$$

$$\chi_{\psi, \nu} \zeta^{\nu^{3/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{3/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \omega_0.$$

$$\chi_{\psi, \nu} \zeta^{\nu^{5/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{3/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \omega_0.$$

$$\chi_{\psi, \nu} \zeta^{\nu} \times \chi_{\psi, \nu} \zeta^{\nu} \times \chi_{\psi, \nu} \zeta^{\nu^{-1/2}} \times \omega_0$$

$$\chi_{\psi, \nu} \zeta^{\nu} \times \chi_{\psi, \nu} \zeta \times \chi_{\psi, \nu} \zeta \times \omega_0$$

$$\chi_{\psi, \nu} \zeta^{\nu^2} \times \chi_{\psi, \nu} \zeta^{\nu} \times \chi_{\psi, \nu} \zeta \times \omega_0$$

$$\chi_{\psi, \nu} \zeta^{\nu} \times \chi_{\psi, \nu} \zeta \times \chi_{\psi, \nu} \zeta^{\nu^{-1}} \times \omega_0$$

$$\chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{-1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{-1/2}} \times \omega_0.$$

$$\chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{-1/2}} \times \omega_0.$$

$$\chi_{\psi, \nu} \zeta^{\nu^{3/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{-1/2}} \times \omega_0.$$

$$\chi_{\psi, \nu} \zeta^{\nu^{1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{-1/2}} \times \chi_{\psi, \nu} \zeta^{\nu^{-3/2}} \times \omega_0.$$

Rank one reducibility

These results are obtained from the work of Muic and Hanzer (2011)

❖ The representation

$$\text{Ind}_{\widetilde{P_1}}^{\widetilde{Sp(3)}} \chi_{V,\psi} \chi_{\nu^s} \otimes \sigma$$

reduces for $s = \{1/2, 3/2, 5/2\}$, depending on the Theta lift of σ .

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reduces for $s = 1/2$.

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❖ The representation

$$\text{Ind}_{\widetilde{P}_3}^{\widetilde{Sp}(3)} \chi_{V,\psi} \chi \nu^s \otimes \omega_0$$

reduces for $s = 1/2$.

❖ The representation

$$\text{Ind}_{\widetilde{P}_2}^{\widetilde{Sp}(3)} \chi_{V,\psi} \chi \nu^s \otimes \sigma$$

reduces for $s = \{0, 1/2, 1, 3/2\}$.

THANK YOU!