

Weighted Low-lying Zeros of L-functions Attached to Siegel Modular Forms

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$$D(f; \Phi) = \sum_{\rho_f} \Phi\left(\frac{\gamma_f}{2\pi} \log c_f\right),$$

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- Since $\Phi(x)$ is of rapid decay, $D(f; \Phi)$ should be regraded as capturing "low-lying" zeros of $L(s, f)$.

- Based on Random Matrix Model prediction, we expect the following:

$$\lim_{Q \rightarrow \infty} \frac{1}{|\mathcal{F}_Q|} \sum_{f \in \mathcal{F}_Q} D(f; \Phi) = \int_{-\infty}^{\infty} \Phi(x) W(G(\mathcal{F}))(x) dx,$$

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- Here $G(\mathcal{F})$ is one of the five classic compact groups: O , $SO(\text{even})$, $SO(\text{odd})$, Sp , U , and $W(G(\mathcal{F}))$ is the distribution determined by low-lying eigenvalues of large random matrices of type $G(\mathcal{F})$.

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- In reality, $\text{supp}(\hat{\Phi})$ must be restricted to certain range. We want this range to be as large as possible.

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- Hecke L -functions of holomorphic cusp forms: \mathbf{O} . (Iwaniec, Luo and Sarnak, 2000)

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- Hasse-Weil L -functions of elliptic curves: **O**. (Young, 2006)
- Spinor (standard) L -functions of genus 2 Siegel modular forms: **O (Sp)**. (Kim, Wakatsuki and Yamauchi, 2020)

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- A Siegel modular form of even weight $k \geq 6$ for Γ_2 is a holomorphic function $F : \mathcal{H}_2 \rightarrow \mathbb{C}$ such that

$$F((AZ + B)(CZ + D)^{-1}) = \det(CZ + D)^k F(Z)$$

for $\begin{pmatrix} A & B \\ C & D \end{pmatrix} \in \Gamma_2$ and $Z \in \mathcal{H}_2$.

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- Denote by $S_k(\Gamma_2)$ the space of Siegel cusp forms of weight k for Γ_2 .
- Each $F \in S_k(\Gamma_2)$ admits a Fourier expansion

$$F(Z) = \sum_{T \in \mathcal{T}} a_F(T) (\det T)^{\frac{k}{2} - \frac{3}{4}} e(\mathrm{Tr}(TZ)),$$

where $\mathcal{T} = \{T = (t_{ij}) \in M_2(\mathbb{R}) : T > 0, t_{11}, t_{22} \in \mathbb{Z}, 2t_{12} = 2t_{21} \in \mathbb{Z}\}$.

Spinor and Standard L-functions

- There is a nice Hecke theory on $S_k(\Gamma_2)$. We thus can choose a Hecke eigenbasis $H_k(\Gamma_2)$. Fix $F \in H_k(\Gamma_2)$ a Hecke eigenform.

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- The spinor L -function $L(s, F; \text{spin})$ is a degree 4 L -function with local parameters at p : $\alpha_p, \beta_p, \alpha_p^{-1}, \beta_p^{-1}$. Its functional equation was proved by Andrianov (1974):

$$\Lambda(s, F; \text{spin}) = \Gamma_{\mathbb{C}}(s + 1/2)\Gamma_{\mathbb{C}}(s + k - 3/2)L(s, F; \text{spin}) = \Lambda(1 - s, F; \text{spin}).$$

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- The standard L -function $L(s, F; \text{std})$ is of degree 5 with local parameters $1, \alpha_p\beta_p, \alpha_p\beta_p^{-1}, \alpha_p^{-1}\beta_p, \alpha_p^{-1}\beta_p^{-1}$. The functional equation was proved by Böcherer (1985):

$$\Lambda(s, F; \text{std}) = \Gamma_{\mathbb{R}}(s)\Gamma_{\mathbb{C}}(s + k - 1)\Gamma_{\mathbb{C}}(s + k - 2)L(s, F; \text{std}) = \Lambda(1 - s, F; \text{std}).$$

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- When F_f is a Saito-Kurokawa lift from $f \in S_{2k-2}(SL_2(\mathbb{Z}))$, we have

$$L(s, F_f; \text{spin}) = \zeta(s + 1/2)\zeta(s - 1/2)L(s, f).$$

Main Results

- Recall that we choose an even Schwartz function $\Phi(x)$. We also choose $c_{F;\text{spin}} = k^2$ and $c_{F;\text{std}} = k^4$, comparable to analytic conductors of $L(s, F; \text{spin})$ and $L(s, F; \text{std})$ resp.

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- With these choices, 1-level density functions $D(F; \Phi; \text{spin})$ and $D(F; \Phi; \text{std})$ are defined.
- For $F \in H_k(\Gamma_2)$ define harmonic weight

$$\omega_F = \frac{\sqrt{\pi}}{4} (4\pi)^{3-2k} \Gamma(k-3/2) \Gamma(k-2) \frac{a_F(l_2)^2}{\|F\|^2}.$$

These weights satisfy

$$\sum_{F \in H_k(\Gamma_2)} \omega_F = 1 + O(e^{-k}).$$

Theorem (Zhao, 2023)

Suppose $\text{supp}(\hat{\Phi}) \subset (-1, 1)$. Assume GRH for $L(s, F; \text{spin})$. Then

$$\lim_{k \rightarrow \infty} \sum_{F \in H_k(\Gamma_2)} \omega_F D(F; \Phi; \text{spin}) = \hat{\Phi}(0) - \frac{\Phi(0)}{2} = \int_{-\infty}^{\infty} \Phi(x) W(\text{Sp})(x) dx.$$

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- Kim, Wakatsuki and Yamauchi (2020) proved the unweighted 1-level density has orthogonal symmetry type.

Main Results

- By extracting terms with $\rho_{F;\text{spin}} = 1/2$ from

$$D(F; \Phi; \text{spin}) = \sum_{\rho_{F;\text{spin}}} \Phi\left(\frac{\gamma_{F;\text{spin}}}{2\pi} \log k^2\right)$$

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Corollary (Zhao, 2023)

Assume GRH for $L(s, F; \text{spin})$. Then we have

$$\liminf_{k \rightarrow \infty} \sum_{\substack{F \in H_k(\Gamma_2) \\ L(1/2, F; \text{spin}) \neq 0}} \omega_F \geq \frac{3}{4}.$$

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- Blomer (2019) used moments method to establish unconditionally

$$\sum_{\substack{F \in H_k(\Gamma_2) \\ L(1/2, F; \text{spin}) \neq 0}} \omega_F \gg \frac{1}{\log k}.$$

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Suppose $\text{supp}(\hat{\Phi}) \subset (-1/4, 1/4)$. Assume GRH for $L(s, F; \text{std})$. Then

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- Kim, Wakatsuki and Yamauchi (2020) found the corresponding unweighted 1-level density has symplectic symmetry type.
- Previously it was not known what the weighted 1-level density symmetry type is. Our result indicates that it remains symplectic.

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Suppose $\text{supp}(\hat{\Phi}) \subset (-5/18, 5/18)$. Assume GRH for $L(s, F; \text{std})$. Then

$$\lim_{K \rightarrow \infty} \frac{2}{K} \sum_{k \sim K} \sum_{F \in H_k(\Gamma_2)} \omega_F D(F; \Phi; \text{std}) = \hat{\Phi}(0) - \frac{\Phi(0)}{2} = \int_{-\infty}^{\infty} \Phi(x) W(\text{Sp})(x) dx.$$

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- Unfortunately, this range of support $(-5/18, 5/18)$ is still not large enough to say anything about non-vanishing of $L(1/2, F; \text{std})$.

- Explicit formula

$$D(F; \Phi; \text{spin}) = \frac{2}{2\pi i} \int_{(2)} \Phi \left(\frac{s - \frac{1}{2}}{2\pi i} \log k^2 \right) \frac{\Lambda'}{\Lambda}(s, F; \text{spin}) ds + (\dots).$$

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- Apply a Petersson-type trace formula (Kitaoka, 1984).
- For averaging over k , apply an integral representation of products of Bessel functions. Then apply the method of Neumann series.